

Non-zero constraints in quantitative coupled physics imaging

Giovanni S. Alberti

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Quantitative Tomographic Imaging – Radon meets Bell and Maxwell

Internal data in quantitative hybrid imaging problems

- ▶ Hybrid conductivity imaging [Widlak, Scherzer, 2012]

$$\begin{cases} -\operatorname{div}(a \nabla u^i) = 0 & \text{in } \Omega, \\ u^i = \varphi_i & \text{on } \partial\Omega. \end{cases}$$

$$u^i(x) \quad \text{or} \quad a(x) \nabla u^i(x) \quad \text{or} \quad a(x) |\nabla u^i|^2(x) \quad \xrightarrow{?} \quad a$$

- ▶ Quantitative thermoacoustic tomography [Bal et al., 2011, Ammari et al., 2013]

$$\begin{cases} \Delta u^i + (\omega^2 + i\omega\sigma) u^i = 0 & \text{in } \Omega, \\ u^i = \varphi_i & \text{on } \partial\Omega. \end{cases}$$

$$\sigma(x) |u^i|^2(x) \quad \xrightarrow{?} \quad \sigma$$

- ▶ MREIT [Seo et al., 2012, Bal and Guo, 2013]

$$\begin{cases} \operatorname{curl} E^i = i\omega H^i & \text{in } \Omega, \\ \operatorname{curl} H^i = -i(\omega\varepsilon + i\sigma) E^i & \text{in } \Omega, \\ E^i \times \nu = \varphi_i \times \nu & \text{on } \partial\Omega. \end{cases}$$

$$H^i(x) \quad \xrightarrow{?} \quad \varepsilon, \sigma$$

The measurements are meaningful at $x \in \Omega$ if at least $u_\omega^i(x) \neq 0$, $\nabla u_\omega^i(x) \neq 0$, ...

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- ▶ Hybrid conductivity imaging [Widlak, Scherzer, 2012], Quantitative PAT

$$\begin{cases} -\operatorname{div}(a \nabla u^i) + \mu u^i = 0 & \text{in } \Omega, \\ u^i = \varphi_i & \text{on } \partial\Omega. \end{cases}$$

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The **measurements** are meaningful at $x \in \Omega$ if at least $u_\omega^i(x) \neq 0$, $\nabla u_\omega^i(x) \neq 0$, ...

Why do non-zero constraints matter?

- ▶ Consider for simplicity the hybrid conductivity problem

$$\begin{cases} -\operatorname{div}(a \nabla u) = 0 & \text{in } \Omega, \\ u = \varphi & \text{on } \partial\Omega. \end{cases}$$

with internal data ∇u and unknown a .

- ▶ With 1 measurement:

$$\nabla a \cdot \nabla u = -a \Delta u \quad \implies \quad \nabla(\log a) \cdot \nabla u = -\Delta u$$

This equation may be solved in a if a is known on $\partial\Omega$ and if

$$\nabla u(x) \neq 0, \quad x \in \Omega.$$

- ▶ With d measurements:

$$\begin{aligned} \nabla(\log a) \cdot (\nabla u^1, \dots, \nabla u^d) &= -(\Delta u^1, \dots, \Delta u^d) \\ \implies \nabla(\log a) &= -(\Delta u^1, \dots, \Delta u^d)(\nabla u^1, \dots, \nabla u^d)^{-1} \end{aligned}$$

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Main question

Is it possible to find suitable illuminations φ_i so that the corresponding solutions u^i satisfy certain non-zero constraints, such as the absence of critical points?

Ideally, we would like to construct the φ_i s a priori, namely independently of the unknown parameters.

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Outline of the talk

- 1 The conductivity equation
- 2 The Helmholtz equation
- 3 The Maxwell's equations



G. S. Alberti and Y. Capdeboscq. Lectures on elliptic methods for hybrid inverse problems. Technical Report 2016-46, SAM, ETH Zürich, 2016.



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1 The conductivity equation

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The Radò-Kneser-Choquet theorem

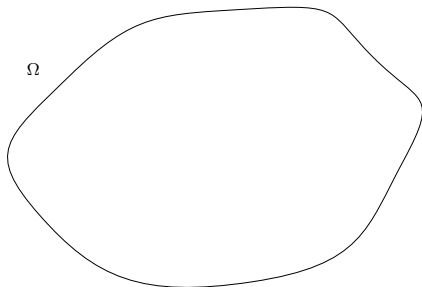
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Let $\Omega \subseteq \mathbb{R}^2$ be a $C^{1,\alpha}$ bounded convex domain and $a \in C^{0,\alpha}(\bar{\Omega}; \mathbb{R}^{2 \times 2})$ be uniformly elliptic. Let $u^i \in H^1(\Omega)$ be the solutions to

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$$\det [\nabla u^1(x) \quad \nabla u^2(x)] \neq 0, \quad x \in \Omega.$$



- ▶ $\det [\nabla u^1(x_0) \quad \nabla u^2(x_0)] = 0$
- ▶ Thus, $\alpha \nabla u^1(x_0) + \beta \nabla u^2(x_0) = 0$
- ▶ Set $v(x) = \alpha u^1(x) + \beta u^2(x)$:
 - ▶ $-\operatorname{div}(a \nabla v) = 0$ in Ω
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- ▶ Thus, v has a saddle point in x_0
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- ▶ But $v(x) = \alpha x_1 + \beta x_2$ on $\partial\Omega$

The Radò-Kneser-Choquet theorem

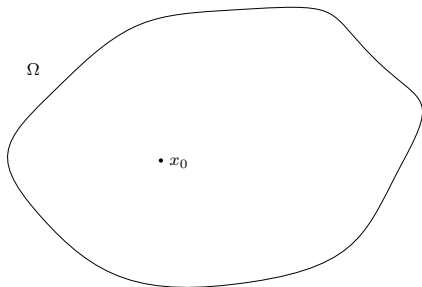
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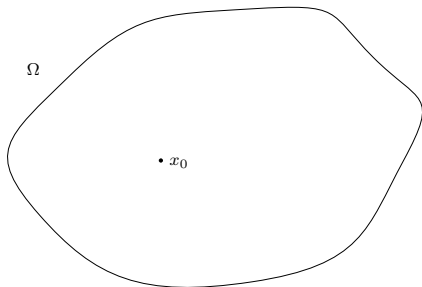
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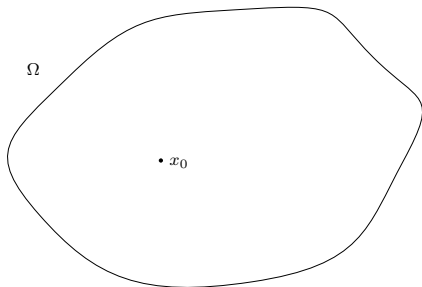
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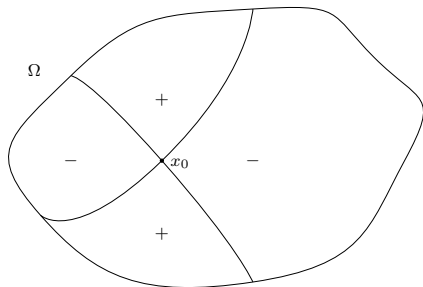
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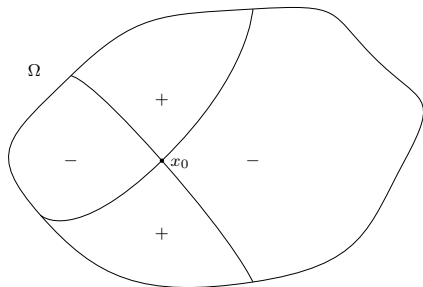
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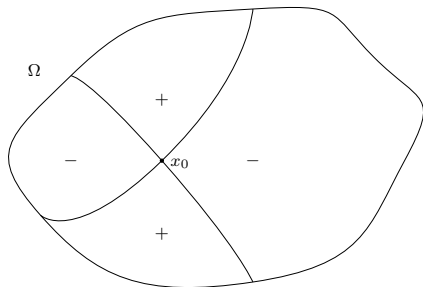
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The failure in three dimensions

$$-\operatorname{div}(a\nabla u^i) = 0 \quad \text{in } \Omega, \quad u^i = \varphi_i \quad \text{on } \partial\Omega.$$

In three dimensions, the above result fails. Several counterexamples:

1. Laugesen 1996: the harmonic case ($a \equiv 1$) for a specific diffeomorphism $\varphi = (\varphi^1, \varphi^2)$
2. Briane et al 2004: the non-constant case (homogenization) for a specific diffeomorphism $\varphi = (\varphi^1, \varphi^2)$
3. Could it be possible to find (φ^1, φ^2) independently of a so that for every $x \in \Omega$

$$\det [\nabla u^1(x) \quad \nabla u^2(x) \quad \nabla u^3(x)] \neq 0?$$

Capdeboscq 2015: No! (by using 2.)

4. What about critical points: can we find φ independently of a so that for every $x \in \Omega$

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The failure in three dimensions

$$-\operatorname{div}(a\nabla u^i) = 0 \quad \text{in } \Omega, \quad u^i = \varphi_i \quad \text{on } \partial\Omega.$$

In three dimensions, the above result fails. Several counterexamples:

1. Laugesen 1996: the harmonic case ($a \equiv 1$) for a specific diffeomorphism $\varphi = (\varphi^1, \varphi^2)$
2. Briane et al 2004: the non-constant case (homogenization) for a specific diffeomorphism $\varphi = (\varphi^1, \varphi^2)$
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Critical points in 3D

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Theorem (GSA, Bal, Di Cristo, ARMA 2017)

Let $\Omega \subseteq \mathbb{R}^3$ be a bounded Lipschitz domain. Take $\varphi \in C(\partial X) \cap H^{\frac{1}{2}}(\partial X)$. There exists a (nonempty open set of) $a \in C^\infty(\overline{X})$ such that the solution $u \in H^1(X)$ to

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has a critical point in Ω , namely $\nabla u(x) = 0$ for some $x \in \Omega$.

Can be extended to deal with:

- ▶ multiple boundary values;
- ▶ multiple critical points (located in arbitrarily small balls);
- ▶ and Neumann boundary conditions.

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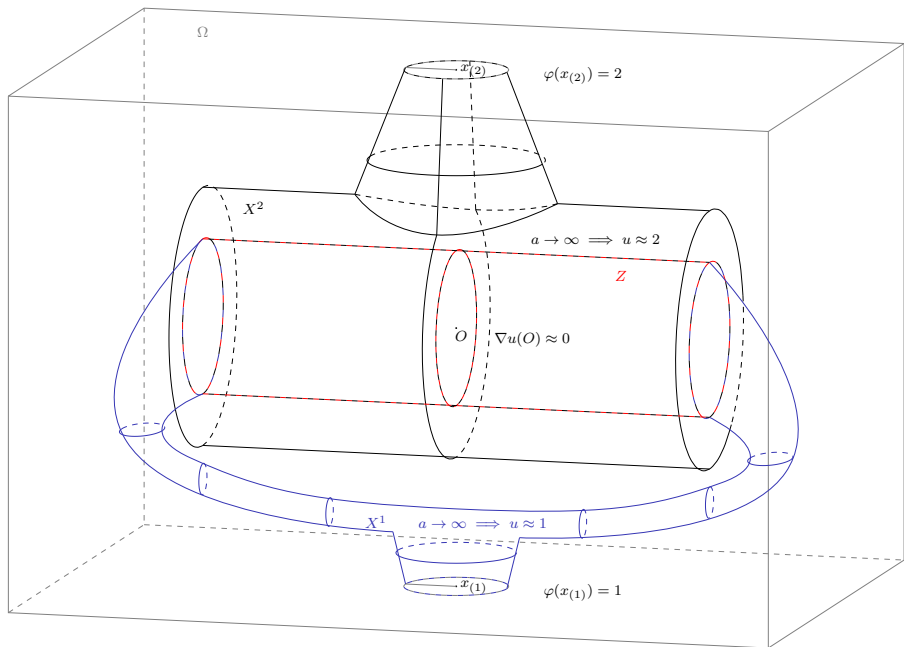
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Alternative approaches

- ▶ **Complex geometrical optics solutions** [Sylvester and Uhlmann, 1987]
 - ▶ $u^{(t)}(x) = e^{tx_m} (\cos(tx_l) + i \sin(tx_l)) (1 + \psi_t)$, $t \gg 1$.
 - ▶ If $t \gg 1$ then $u^{(t)}(x) \approx e^{tx_m} (\cos(tx_l) + i \sin(tx_l))$ in C^1 [Bal and Uhlmann, 2010]
 - ▶ The traces on the boundary of these solutions give the required φ_i s
 - ▶ Need smooth coefficients, construction depends on coefficients.
 - ▶ Only for isotropic coefficients
- ▶ **Runge approximation** [Lax 1956, Bal and Uhlmann 2013]
 - ▶ There exist solutions that are locally closed to the solutions of the constant coefficient PDE.
 - ▶ Based on unique continuation, non constructive.
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- ▶ **Stability results without the constraints**
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Outline of the talk

1 The conductivity equation

2 The Helmholtz equation

3 The Maxwell's equations

The Helmholtz equation

- ▶ We now consider the Helmholtz equation

$$\begin{cases} \Delta u_\omega^i + (\omega^2 \varepsilon + i\omega\sigma) u_\omega^i = 0 & \text{in } \Omega, \\ u_\omega^i = \varphi_i & \text{on } \partial\Omega. \end{cases}$$

where $\Omega \subseteq \mathbb{R}^d$, $d = 2, 3$, $\varepsilon, \sigma \in L^\infty(\Omega)$, $\sigma, \varepsilon \leq \Lambda$, $\varepsilon > \Lambda^{-1}$.

- ▶ We are interested in the constraints:

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2. $|\det [\nabla u_\omega^2 \quad \dots \quad \nabla u_\omega^{d+1}]|(x) > 0$ (Jacobian)
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- ▶ CGO solutions and the Runge approximation may be used also in this case, but the corresponding boundary conditions φ_i are not explicitly constructed.
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Multi-Frequency Approach: main result

$K^{(n)}$: uniform partition of $\mathcal{A} = [K_{min}, K_{max}]$ with n points



Theorem (GSA, IP 2013 & CPDE 2015)

There exist $C > 0$ and $n \in \mathbb{N}^*$ depending only on Ω , Λ and \mathcal{A} such that the following is true. Take

$$\varphi_1 = 1, \quad \varphi_2 = x_1, \quad \dots \quad \varphi_{d+1} = x_d.$$

There exists an open cover

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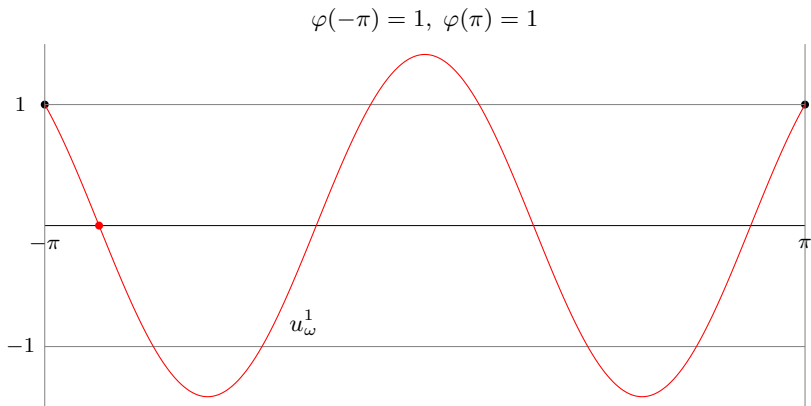
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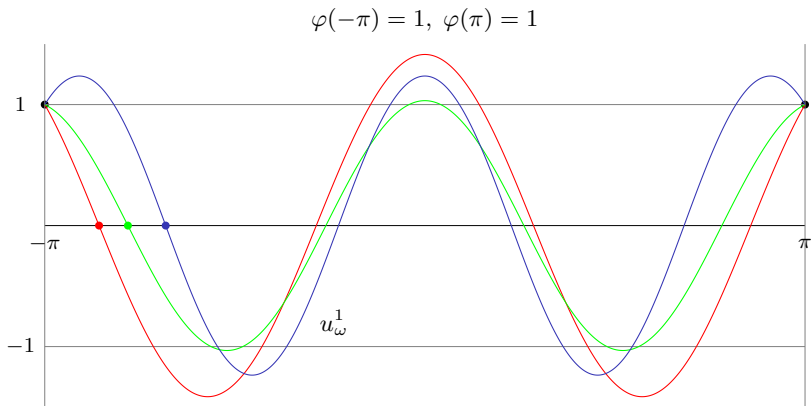
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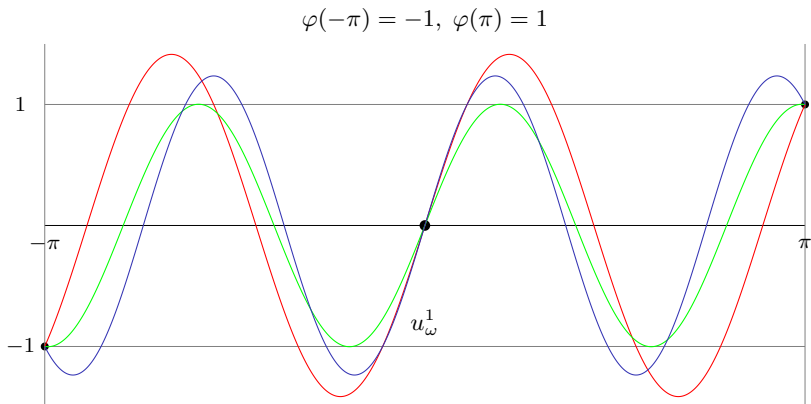


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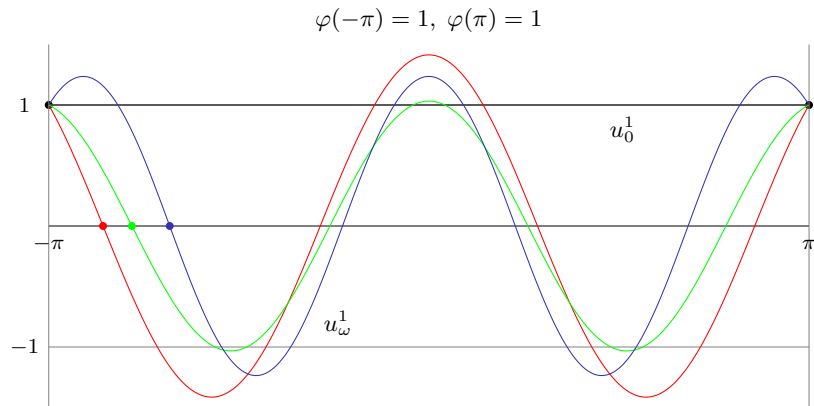
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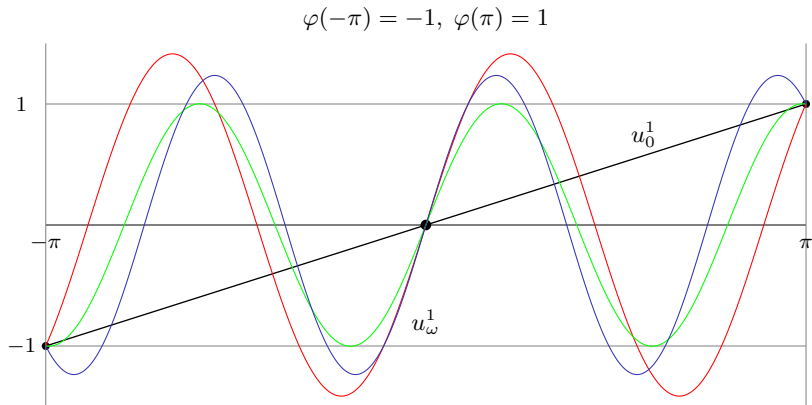
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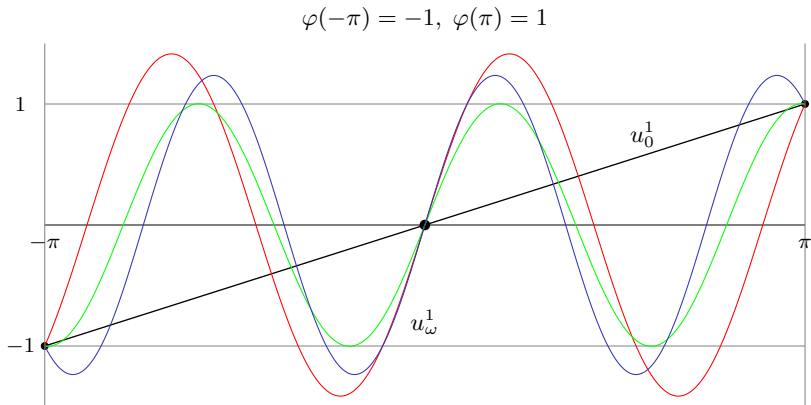
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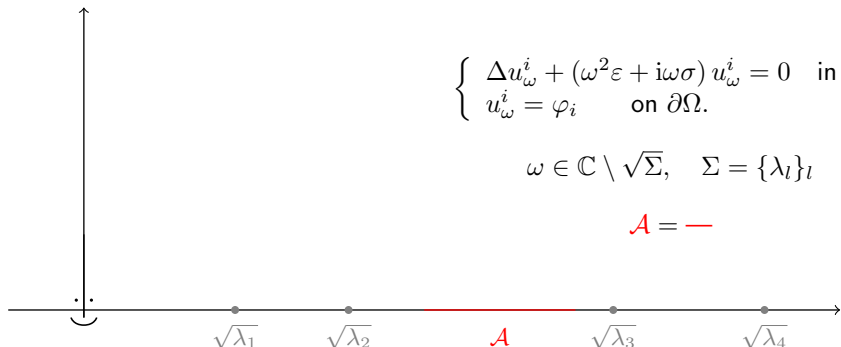
How to pass from 0 to ω ?

Lemma

The map $\mathbb{C} \setminus \sqrt{\Sigma} \rightarrow C^1(\overline{\Omega})$, $\omega \mapsto u_\omega^i$ is holomorphic.

- ▶ The set $Z_x = \{\omega \in \mathcal{A} : u_\omega^1(x) = 0\}$ is finite (consider 1. for simplicity)
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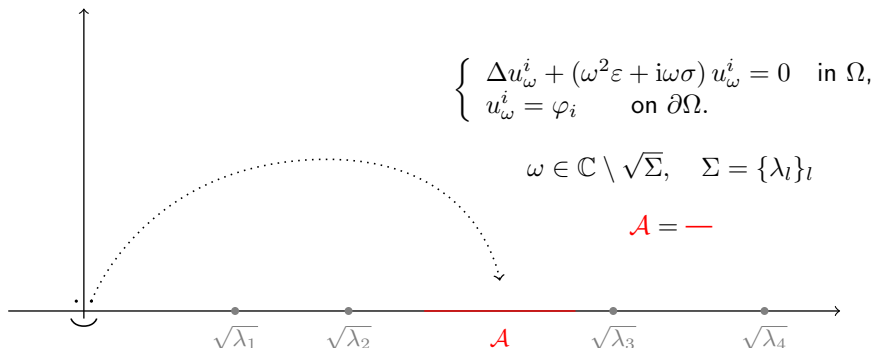


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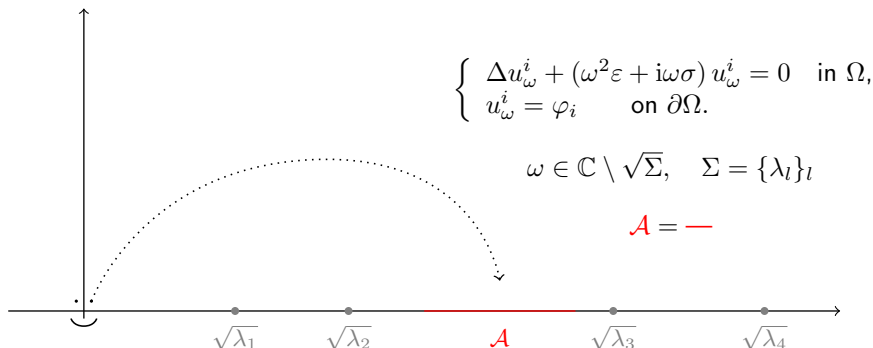


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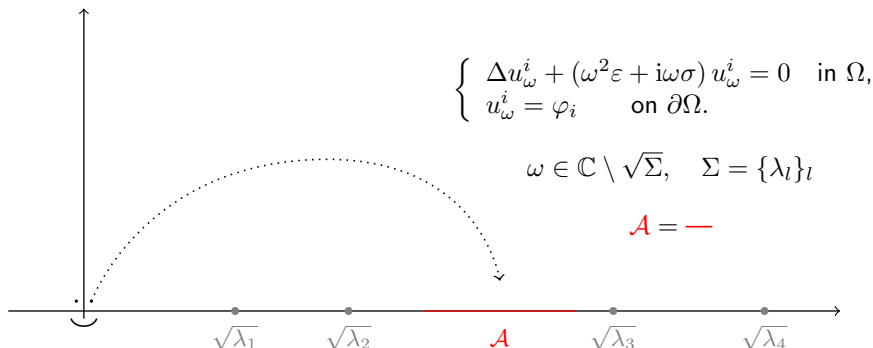


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How many frequencies are needed?

Theorem (GSA and Capdeboscq, CM 2016)

Take $\varphi = 1$. Assume that σ and ε are *real analytic*. The set

$$\left\{ (\omega_1, \dots, \omega_{d+1}) \in \mathcal{A}^{d+1} : \min_{\overline{\Omega}} (|u_{\omega_1}^\varphi| + \dots + |u_{\omega_{d+1}}^\varphi|) > 0 \right\}$$

is open and dense in \mathcal{A}^{d+1} .

In other words, (almost any) $d + 1$ frequencies are ok.

Proof.

- ▶ Classical elliptic regularity theory implies that u_ω^φ is real analytic
- ▶ The set $X = \{x \in \Omega : |u_{\omega_1}^\varphi| = \dots = |u_{\omega_l}^\varphi| = 0\}$ is an analytic variety
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is open and dense in \mathcal{A}^{d+1} .

In other words, (almost any) $d + 1$ frequencies are ok.

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- ▶ Ammari et al. (2016) have successfully adapted this method to

$$\operatorname{div}((\omega\varepsilon + i\sigma)\nabla u_\omega^i) = 0.$$

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Theorem (GSA, ARMA 2016)

Suppose $a, \varepsilon \in C^2(\mathbb{R}^3)$ and $\sigma = 0$. For a generic C^2 bounded domain Ω and a generic $\varphi \in C^1(\overline{\Omega})$ there exists a finite $K \subseteq \mathcal{A}$ such that

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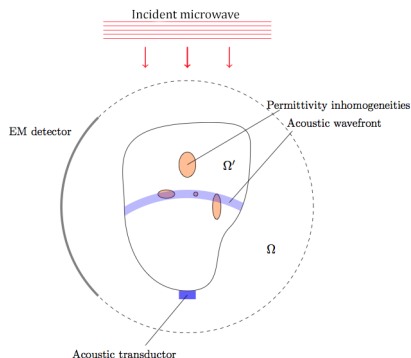
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Acousto-electromagnetic tomography (Ammari et al., 2012)



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$$D\psi_\omega[\varepsilon](\rho) \mapsto \rho$$

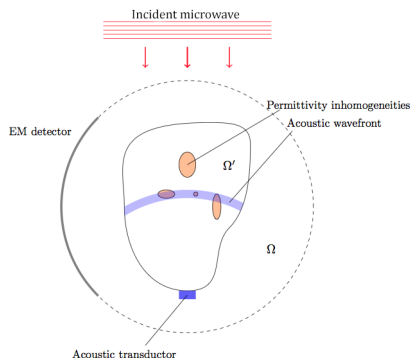
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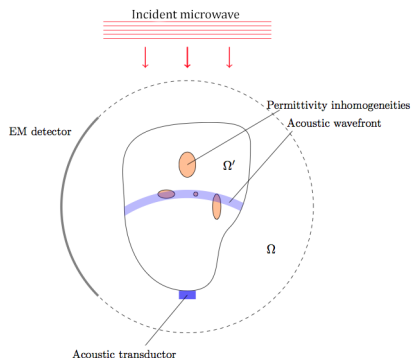
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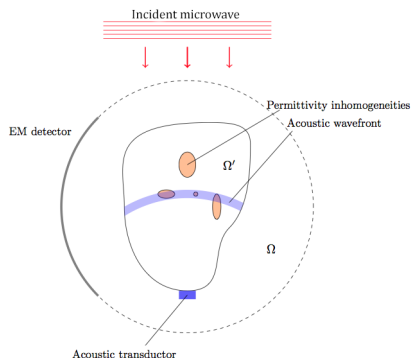
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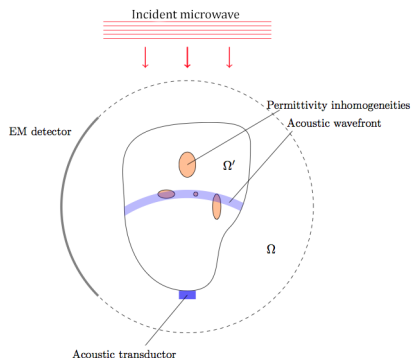
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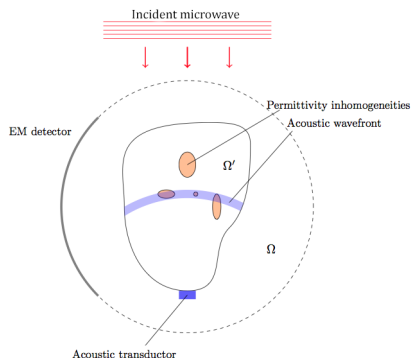
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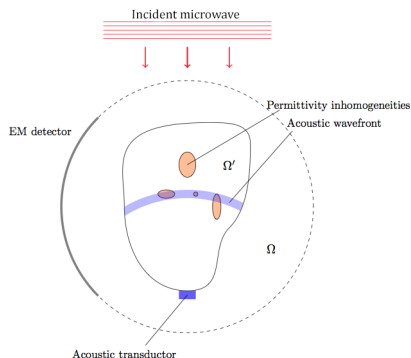
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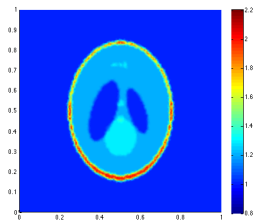
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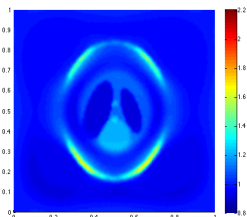
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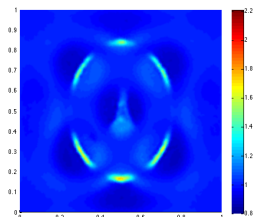
Numerical experiments



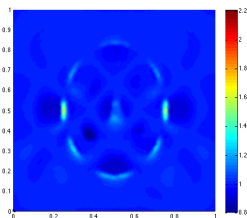
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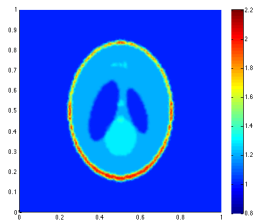
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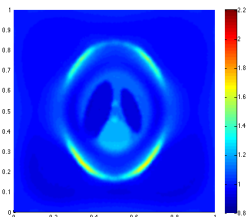
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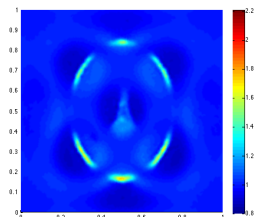
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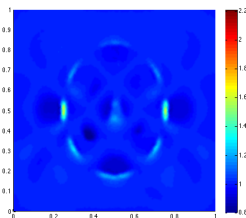
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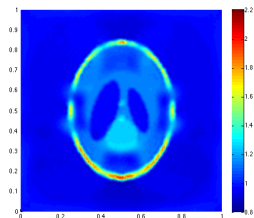
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Outline of the talk

1 The conductivity equation

2 The Helmholtz equation

3 The Maxwell's equations

Constraints for Maxwell's equations

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$$\begin{cases} \operatorname{curl} E^i = i\omega H^i & \text{in } \Omega, \\ \operatorname{curl} H^i = -i(\omega\varepsilon + i\sigma)E^i & \text{in } \Omega, \\ E^i \times \nu = \varphi_i \times \nu & \text{on } \partial\Omega. \end{cases}$$

$$H^i(x) \xrightarrow{?} \varepsilon, \sigma$$

- ▶ The relevant constraint in this case is

$$|\det [E^1(x) \quad E^2(x) \quad E^3(x)]| > 0.$$

- ▶ CGO solutions may be used [Chen Yang, IP 2013].
- ▶ The multi-frequency method discussed above works as well [GSA, JDE 2015].
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Theorem (GSA, 2016)

Assume that

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and that are uniformly elliptic (constant Λ). Take $J_e, J_m \in C^{0,\alpha}(\bar{\Omega}; \mathbb{C}^3)$ and $G \in C^{1,\alpha}(\text{curl}, \Omega)$. Let $(E, H) \in H(\text{curl}, \Omega)^2$ be a weak solution of

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Take-home message: regularity for Maxwell is exactly as in the elliptic case!

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